

Gibbsian germ-grain models

David Dereudre,
Laboratoire de Mathématiques Paul Painlevé,
University Lille 1, France

Workshop, Point Processes and Random Geometry, 7th
October 2013, Bochum

- 1 Germ-grain models
- 2 Infinite volume Gibbsian germ-grain models
- 3 A percolation result
- 4 Parametric estimation for Quermass models

1 Germ-grain models

State space

- $\mathcal{E} = \mathbb{R}^2 \times \mathbb{R}^+$. We denote by $X = (x, R)$ an element in \mathcal{E} .

State space

- $\mathcal{E} = \mathbb{R}^2 \times \mathbb{R}^+$. We denote by $X = (x, R)$ an element in \mathcal{E} .
- $\mathcal{M}(\mathcal{E})$ the space of locally finite configurations in \mathcal{E} . We denote by γ a configuration in $\mathcal{M}(\mathcal{E})$ and by $\bar{\gamma}$ its associated germ-grain set

$$\bar{\gamma} = \bigcup_{(x,R) \in \gamma} B(x, R).$$

State space

- $\mathcal{E} = \mathbb{R}^2 \times \mathbb{R}^+$. We denote by $X = (x, R)$ an element in \mathcal{E} .
- $\mathcal{M}(\mathcal{E})$ the space of locally finite configurations in \mathcal{E} . We denote by γ a configuration in $\mathcal{M}(\mathcal{E})$ and by $\bar{\gamma}$ its associated germ-grain set

$$\bar{\gamma} = \bigcup_{(x,R) \in \gamma} B(x, R).$$

- A Point Process Γ is a random variable in $\mathcal{M}(\mathcal{E})$ and $\bar{\Gamma}$ is its associated germ grain model.

State space

- $\mathcal{E} = \mathbb{R}^2 \times \mathbb{R}^+$. We denote by $X = (x, R)$ an element in \mathcal{E} .
- $\mathcal{M}(\mathcal{E})$ the space of locally finite configurations in \mathcal{E} . We denote by γ a configuration in $\mathcal{M}(\mathcal{E})$ and by $\bar{\gamma}$ its associated germ-grain set

$$\bar{\gamma} = \bigcup_{(x,R) \in \gamma} B(x, R).$$

- A Point Process Γ is a random variable in $\mathcal{M}(\mathcal{E})$ and $\bar{\Gamma}$ is its associated germ grain model.
- First possible extension : The grains can be more complicated (convex sets, paths, etc)

State space

- $\mathcal{E} = \mathbb{R}^2 \times \mathbb{R}^+$. We denote by $X = (x, R)$ an element in \mathcal{E} .
- $\mathcal{M}(\mathcal{E})$ the space of locally finite configurations in \mathcal{E} . We denote by γ a configuration in $\mathcal{M}(\mathcal{E})$ and by $\bar{\gamma}$ its associated germ-grain set

$$\bar{\gamma} = \bigcup_{(x,R) \in \gamma} B(x, R).$$

- A Point Process Γ is a random variable in $\mathcal{M}(\mathcal{E})$ and $\bar{\Gamma}$ is its associated germ grain model.
- First possible extension : The grains can be more complicated (convex sets, paths, etc)
- Second possible extension : The grains can be marked by a type (color). \mathcal{E} becomes for example $\mathbb{R}^2 \times \mathbb{R}^+ \times \{1, \dots, K\}$.

Boolean model

- Q is a probability measure on \mathbb{R}^+ .

Boolean model

- Q is a probability measure on \mathbb{R}^+ .
- π^z is the law of a Poisson Point Process on \mathcal{E} with intensity $z\lambda \otimes Q$.

Boolean model

- Q is a probability measure on \mathbb{R}^+ .
- π^z is the law of a Poisson Point Process on \mathcal{E} with intensity $z\lambda \otimes Q$.
- A boolean model with intensity z is a random set $\bar{\Gamma}$ associated to a Poisson Point Process Γ with intensity z . We denote by P^z the law of this Boolean model in the space of closed sets.

Boolean model

- Q is a probability measure on \mathbb{R}^+ .
- π^z is the law of a Poisson Point Process on \mathcal{E} with intensity $z\lambda \otimes Q$.
- A boolean model with intensity z is a random set $\bar{\Gamma}$ associated to a Poisson Point Process Γ with intensity z . We denote by P^z the law of this Boolean model in the space of closed sets.
- On a bounded window Λ : $\pi_\Lambda^z, P_\Lambda^z$



Gibbsian modifications

On a finite window Λ .

$$P^{z,H} = \frac{1}{Z_\Lambda} e^{-H} P_\Lambda^z.$$

with H an Hamiltonian which depends on $\bar{\gamma}_\Lambda$: a function from the space of finite union of balls to $\mathbb{R} \cup \{+\infty\}$ such that

$$0 < Z_\Lambda := \int e^{-H(\bar{\gamma}_\Lambda)} P_\Lambda^z(d\gamma_\Lambda) < +\infty.$$

Gibbsian modifications

On a finite window Λ .

$$P^{z,H} = \frac{1}{Z_\Lambda} e^{-H} P_\Lambda^z.$$

with H an Hamiltonian which depends on $\bar{\gamma}_\Lambda$: a function from the space of finite union of balls to $\mathbb{R} \cup \{+\infty\}$ such that

$$0 < Z_\Lambda := \int e^{-H(\bar{\gamma}_\Lambda)} P_\Lambda^z(d\gamma_\Lambda) < +\infty.$$

- $Z_\Lambda > 0$: $H(\emptyset) = 0$.
- $Z_\Lambda < +\infty$: Stability

$$H(\bar{\gamma}_\Lambda) \geq -B \text{Card}(\gamma_\Lambda).$$

Examples of Hamiltonian

- **Quermass interaction** (Likos, Mecke and Wagner 95-Baddeley, Van Lieshout 99)

$$H(\bar{\gamma}) = \sum_{i=0}^d \theta_i W_i.$$

with θ_i in \mathbb{R} and W_i the i th Minkowski functional.
(area, perimeter and Euler-Poincaré characteristic in dimension 2)

Examples of Hamiltonian

- **Quermass interaction** (Likos, Mecke and Wagner 95-Baddeley, Van Lieshout 99)

$$H(\bar{\gamma}) = \sum_{i=0}^d \theta_i W_i.$$

with θ_i in \mathbb{R} and W_i the i th Minkowski functional. (area, perimeter and Euler-Poincaré characteristic in dimension 2)

- **Continuum random cluster model** (Georgii, Häggström, Maes 99)

$$H(\bar{\gamma}) = \theta N_{cc}(\bar{\gamma}), \quad \theta \in \mathbb{R}$$

Examples of Hamiltonian

- **Quermass interaction** (Likos, Mecke and Wagner 95-Baddeley, Van Lieshout 99)

$$H(\bar{\gamma}) = \sum_{i=0}^d \theta_i W_i.$$

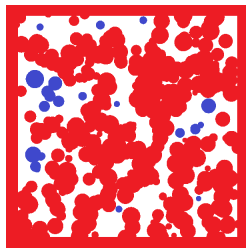
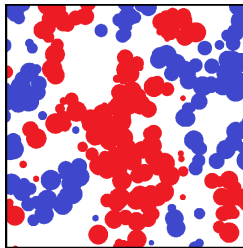
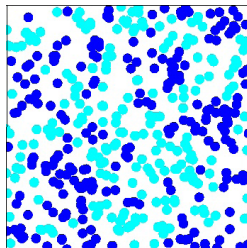
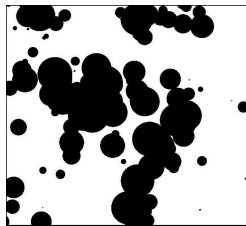
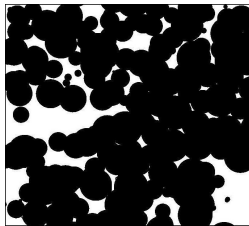
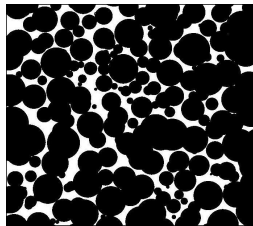
with θ_i in \mathbb{R} and W_i the i th Minkowski functional. (area, perimeter and Euler-Poincaré characteristic in dimension 2)

- **Continuum random cluster model** (Georgii, Häggström, Maes 99)

$$H(\bar{\gamma}) = \theta N_{cc}(\bar{\gamma}), \quad \theta \in \mathbb{R}$$

- **Multi type Widom Rowlinson model** (Widom and Rowlinson 70)

Inhibition model : non overlapping balls with different type.



2 Infinite volume Gibbsian germ-grain models

Why the infinite volume?

Motivations :

- Stationary model without boundary conditions
- Macroscopic quantities (mean value, percolation, conductivity, permeability)
- Phase transition via the non uniqueness of the Gibbs measures
- Asymptotic properties of statistical estimations (consistency, normality, etc)

Why the infinite volume?

Motivations :

- Stationary model without boundary conditions
- Macroscopic quantities (mean value, percolation, conductivity, permeability)
- Phase transition via the non uniqueness of the Gibbs measures
- Asymptotic properties of statistical estimations (consistency, normality, etc)

Issues :

- The infinite volume Hamiltonian is senseless
- Definition of local Hamiltonian
- Equilibrium equations via DLR equations

The Local Hamiltonian

$$H_{\Lambda}(\tilde{\gamma}) = \lim_{\Delta \rightarrow \mathbb{R}^2} H(\tilde{\gamma}_{\Delta}) - H(\tilde{\gamma}_{\Delta \setminus \Lambda})$$

The Local Hamiltonian

$$H_{\Lambda}(\tilde{\gamma}) = \lim_{\Delta \rightarrow \mathbb{R}^2} H(\tilde{\gamma}_{\Delta}) - H(\tilde{\gamma}_{\Delta \setminus \Lambda})$$

In some cases : a space Ω^* of tempered configurations is needed.

The Local Hamiltonian

$$H_{\Lambda}(\bar{\gamma}) = \lim_{\Delta \rightarrow \mathbb{R}^2} H(\bar{\gamma}_{\Delta}) - H(\bar{\gamma}_{\Delta \setminus \Lambda})$$

In some cases : a space Ω^* of tempered configurations is needed.

Definition (Gibbs measures)

An infinite volume germ-grain model P is a Gibbs measure for the Hamiltonian (H_{Λ}) if $P(\Omega^*) = 1$ and if the law of γ_{Λ} given γ_{Λ^c} is absolutely continuous with respect to the Poisson law π_{Λ}^z with the density

$$\frac{1}{Z_{\Lambda}(\gamma_{\Lambda^c})} e^{-H_{\Lambda}(\bar{\gamma})}$$

The Local Hamiltonian

$$H_{\Lambda}(\bar{\gamma}) = \lim_{\Delta \rightarrow \mathbb{R}^2} H(\bar{\gamma}_{\Delta}) - H(\bar{\gamma}_{\Delta \setminus \Lambda})$$

In some cases : a space Ω^* of tempered configurations is needed.

Definition (Gibbs measures)

An infinite volume germ-grain model P is a Gibbs measure for the Hamiltonian (H_{Λ}) if $P(\Omega^*) = 1$ and if the law of γ_{Λ} given γ_{Λ^c} is absolutely continuous with respect to the Poisson law π_{Λ}^z with the density

$$\frac{1}{Z_{\Lambda}(\gamma_{\Lambda^c})} e^{-H_{\Lambda}(\bar{\gamma})}$$

Questions : Existence, uniqueness, non-uniqueness (phase transition).

Some results

- **Widom Rowlinson exclusion model :**
(Widom-Rowlinson 70, Chayes Kotecky 95) The law of radii Q is concentrated on a singleton R_0 : Existence and phase transition for large z large enough.

Some results

- **Widom Rowlinson exclusion model :**
(Widom-Rowlinson 70, Chayes Kotecky 95) The law of radii Q is concentrated on a singleton R_0 : Existence and phase transition for large z large enough.
- **Quermass model :** (Der. 2009) Existence under the assumption

$$\forall \theta > 0, \quad \int e^{\theta R^2} Q(dR) < +\infty.$$

Some results

- **Widom Rowlinson exclusion model :**
(Widom-Rowlinson 70, Chayes Kotecky 95) The law of radii Q is concentrated on a singleton R_0 : Existence and phase transition for large z large enough.
- **Quermass model :** (Der. 2009) Existence under the assumption

$$\forall \theta > 0, \quad \int e^{\theta R^2} Q(dR) < +\infty.$$

- **Continuum random cluster model :** work in progress with my Phd Student, Pierre Houdebert. Existence and phase transition.

Stability of Quermass Model in \mathbb{R}^2

Is e^{-H} intégrable under π_Λ ?

Proposition (KVB99, MH08)

If

$$\int_{\mathbb{R}^+} e^{-\theta_1 \pi R^2 - 2\pi \theta_2 R} Q(dr) < +\infty,$$

then

$$\int e^{-H(\gamma)} \pi_\Lambda(d\gamma) < \infty.$$

Proof :

For one ball : $\int e^{-\theta_1 \mathcal{A}(B(x,R)) - \theta_2 \mathcal{L}(B(x,R))} Q(dR)$

Lemme (KVB99)

Let n ($n \geq 3$) balls be in the plane. Then the number of holes is lower than $2n - 5$.

H is stable : $H(\gamma) \geq -K \text{Card}(\gamma)$.

Tempered configurations for Quermass Model

$$\Omega_{K,K'} = \left\{ \gamma \text{ tq} \begin{array}{l} -i) \sup_{n \in \mathbb{N}^*} \frac{1}{\pi n^2} \sum_{(x,R) \in \gamma_{B(0,n)}} (1 + R^2) \leq K \\ -ii) \forall n \in \mathbb{N}^*, \sup_{(x,R) \in \gamma_{B(0,n)}} R \leq \frac{1}{2}n + K' \end{array} \right\}.$$

$$\Omega^* = \bigcup_{K \geq 2, K' \geq 2} M_{K,K'}(\mathcal{E}).$$

Tempered configurations for Quermass Model

$$\Omega_{K,K'} = \left\{ \gamma \text{ tq} \begin{array}{l} -i) \sup_{n \in \mathbb{N}^*} \frac{1}{\pi n^2} \sum_{(x,R) \in \gamma_{B(0,n)}} (1 + R^2) \leq K \\ -ii) \forall n \in \mathbb{N}^*, \sup_{(x,R) \in \gamma_{B(0,n)}} R \leq \frac{1}{2}n + K' \end{array} \right\}.$$

$$\Omega^* = \bigcup_{K \geq 2, K' \geq 2} M_{K,K'}(\mathcal{E}).$$

The local Hamiltonian

$$H_\Lambda(\bar{\gamma}) = \lim_{\Delta \rightarrow \mathbb{R}^2} H(\bar{\gamma}_\Delta) - H(\bar{\gamma}_{\Delta \setminus \Lambda})$$

is well defined for tempered configurations.

Tempered configurations for Quermass Model

$$\Omega_{K,K'} = \left\{ \gamma \text{ tq} \begin{array}{l} -i) \sup_{n \in \mathbb{N}^*} \frac{1}{\pi n^2} \sum_{(x,R) \in \gamma_{B(0,n)}} (1 + R^2) \leq K \\ -ii) \forall n \in \mathbb{N}^*, \sup_{(x,R) \in \gamma_{B(0,n)}} R \leq \frac{1}{2}n + K' \end{array} \right\}.$$

$$\Omega^* = \bigcup_{K \geq 2, K' \geq 2} M_{K,K'}(\mathcal{E}).$$

The local Hamiltonian

$$H_\Lambda(\bar{\gamma}) = \lim_{\Delta \rightarrow \mathbb{R}^2} H(\bar{\gamma}_\Delta) - H(\bar{\gamma}_{\Delta \setminus \Lambda})$$

is well defined for tempered configurations.

The cluster points in the construction of Gibbs measures by "entropy bounds" is tempered.

3 A percolation result

A Result

" $\bar{\gamma}$ percolates" means "there exists an unbounded connected component in $\bar{\gamma}$ ".

A Result

" $\bar{\gamma}$ percolates" means "there exists an unbounded connected component in $\bar{\gamma}$ ".

Theorem (Couplier, Der.. 2012)

We assume that $Q([R_0, R_1]) = 1$ with $(R_0 > 0$ and $R_1 < \infty)$, then for any coefficients $\theta_1, \theta_2, \theta_3$ in \mathbb{R} , there exists $z^ > 0$ such that for any $z > z^*$ and any Quermass process P for parameters $z, \theta_1, \theta_2, \theta_3$,*

$$P(\bar{\gamma} \text{ percolates}) = 1,$$

A Result

" $\bar{\gamma}$ percolates" means "there exists an unbounded connected component in $\bar{\gamma}$ ".

Theorem (Coupier, Der.. 2012)

We assume that $Q([R_0, R_1]) = 1$ with $(R_0 > 0$ and $R_1 < \infty)$, then for any coefficients $\theta_1, \theta_2, \theta_3$ in \mathbb{R} , there exists $z^ > 0$ such that for any $z > z^*$ and any Quermass process P for parameters $z, \theta_1, \theta_2, \theta_3$,*

$$P(\bar{\gamma} \text{ percolates}) = 1,$$

Remark : There exists Quermass process P such that

$$0 < P(\bar{\gamma} \text{ percolates}) < 1.$$

A Result

" $\bar{\gamma}$ percolates" means "there exists an unbounded connected component in $\bar{\gamma}$ ".

Theorem (Coupier, Der.. 2012)

We assume that $Q([R_0, R_1]) = 1$ with $(R_0 > 0$ and $R_1 < \infty)$, then for any coefficients $\theta_1, \theta_2, \theta_3$ in \mathbb{R} , there exists $z^ > 0$ such that for any $z > z^*$ and any Quermass process P for parameters $z, \theta_1, \theta_2, \theta_3$,*

$$P(\bar{\gamma} \text{ percolates}) = 1,$$

Remark : There exists Quermass process P such that

$$0 < P(\bar{\gamma} \text{ percolates}) < 1.$$

Main issue : when $\theta_3 \neq 0$, it is impossible to obtain a stochastic minoration of P by Poisson processes

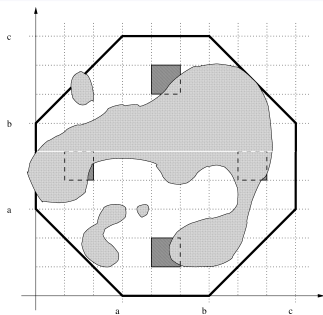
$$\text{For all } z' > 0, \quad \pi_{\Lambda}^{z'} \not\leq P_{\Lambda}.$$

The connection Lemma

\mathcal{D} = the diamond box

\mathcal{D} is open for $\bar{\gamma}$ if

- $\bar{\gamma} \cap B_N \neq \emptyset$
- the same for B_E, B_W, B_S
- B_N, B_E, B_W, B_S are connected via $\bar{\gamma}_{\mathcal{D}}$

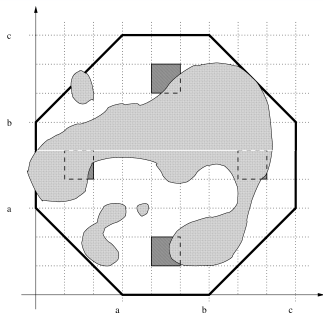


The connection Lemma

\mathcal{D} = the diamond box

\mathcal{D} is open for $\bar{\gamma}$ if

- $\bar{\gamma} \cap B_N \neq \emptyset$
- the same for B_E, B_W, B_S
- B_N, B_E, B_W, B_S are connected via $\bar{\gamma}_{\mathcal{D}}$



Lemma (Connection Lemma)

There exists $C > 0$ (depending on θ_1, θ_2 and θ_3) such that for any $z > 0$ and any Quermass process P

$$\inf_{\gamma_{\Lambda^c}} P(\mathcal{D} \text{ is open} \mid \gamma_{\Lambda^c}) \geq 1 - \frac{C}{z}.$$

Classical Bernoulli domination

Let (V, E) be an undirected graph with uniformly bounded degrees and ξ a random variable in $\{0, 1\}^V$

Lemme (Liggett et al. 97)

Let $p \in [0, 1]$. Assume that for all $x \in V$,

$$P(\xi_x = 1 \mid \xi_y : \{x, y\} \notin E) \geq p \text{ a.s.}$$

Then the law of $\{\xi_x, x \in V\}$ dominates stochastically a product $\otimes_{x \in V} B_x$ of Bernoulli laws with parameter $f(p)$, with $\lim_{p \rightarrow 1} f(p) = 1$.

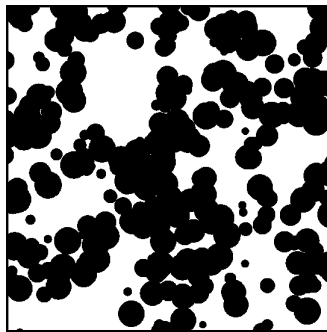
Representation of the multi-type Quermass model on Λ

For $K = 2$:

A one-type Quermass model P_Λ on Λ with density $2^{N_{cc}(\bar{\gamma})}$:

$$Q_\Lambda(d\gamma) = \frac{1}{Z_\Lambda} 2^{N_{cc}(\bar{\gamma})} P_\Lambda(d\gamma).$$

Example with $\theta_1 = -0.2$, $\theta_2 = 0.3$ and $\theta_3 = 0$:



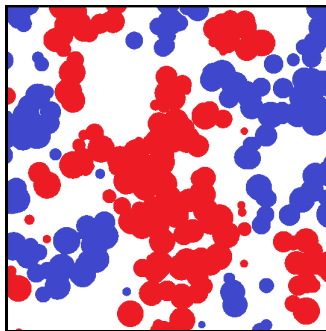
Representation of the multi-type Quermass model on Λ

For $K = 2$:

A one-type Quermass model P_Λ on Λ with density $2^{N_{cc}(\bar{\gamma})}$:

$$Q_\Lambda(d\gamma) = \frac{1}{Z_\Lambda} 2^{N_{cc}(\bar{\gamma})} P_\Lambda(d\gamma).$$

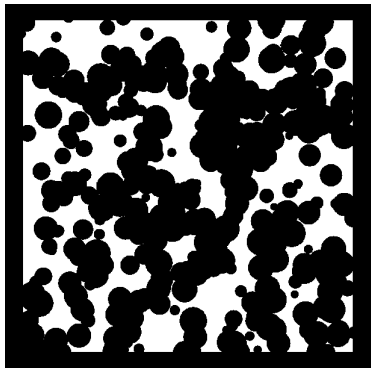
Example with $\theta_1 = -0.2$, $\theta_2 = 0.3$ and $\theta_3 = 0$:



In colouring independently the connected components, we obtain a 2-type Quermass model on Λ for the same parameters.

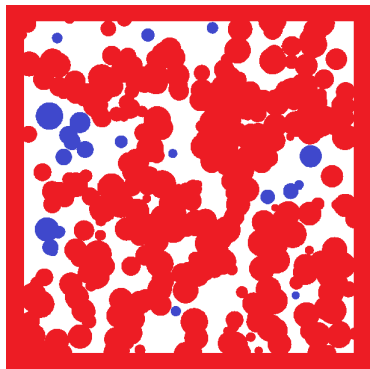
The phase transition proof

Let $\frac{1}{Z_\Lambda} 2^{N_{cc}} P_\Lambda^z(\cdot | \gamma_{\Lambda^c})$ be a modified one-type Quermass process with a full boundary condition



The phase transition proof

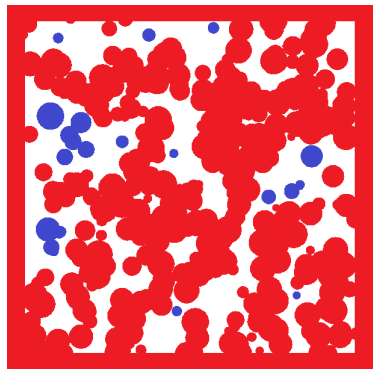
Let $\frac{1}{Z_\Lambda} 2^{N_{cc}} P_\Lambda^z(\cdot | \gamma_{\Lambda^c})$ be a modified one-type Quermass process with a full boundary condition



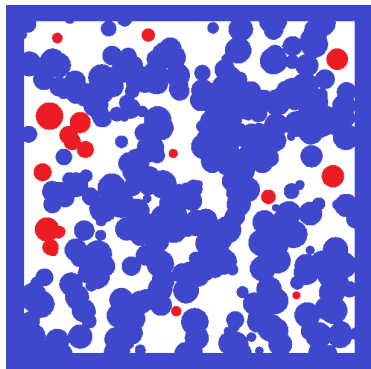
2-type Quermass Process in Λ
with red boundary condition

The phase transition proof

Let $\frac{1}{Z_\Lambda} 2^{N_{cc}} P_\Lambda^z(\cdot | \gamma_{\Lambda^c})$ be a modified one-type Quermass process with a full boundary condition



2-type Quermass Process in Λ
with red boundary condition



2-type Quermass Process in Λ
with blue boundary condition

The phase transition proof

- When Λ goes to \mathbb{R}^2 , the 2-type Quermass process in Λ with red boundary condition goes to a 2-type Quermass process in \mathbb{R}^2 with the red particle density bigger than the blue particle density (if percolation occurs).

The phase transition proof

- When Λ goes to \mathbb{R}^2 , the 2-type Quermass process in Λ with red boundary condition goes to a 2-type Quermass process in \mathbb{R}^2 with the red particle density bigger than the blue particle density (if percolation occurs).
- Conversely for the 2-type Quermass process in Λ with blue boundary condition.

The phase transition proof

- When Λ goes to \mathbb{R}^2 , the 2-type Quermass process in Λ with red boundary condition goes to a 2-type Quermass process in \mathbb{R}^2 with the red particle density bigger than the blue particle density (if percolation occurs).
- Conversely for the 2-type Quermass process in Λ with blue boundary condition.
- We build two different 2-type Quermass processes in \mathbb{R}^2 .

4 Parametric estimation for Quermass models

MLE and MPLE procedures

- P a Quermass Model for $\Theta^* = (z^*, \theta_1^*, \theta_2^*, \theta_3^*)$.
- h^{Θ^*} and h^Θ are the local energies for Θ^* and Θ .
- γ a realization of P .

MLE and MPLE procedures

- P a Quermass Model for $\Theta^* = (z^*, \theta_1^*, \theta_2^*, \theta_3^*)$.
- h^{Θ^*} and h^Θ are the local energies for Θ^* and Θ .
- γ a realization of P .
- MLE** : In Λ , the MLE is defined by

$$\hat{\Theta}_\Lambda = \operatorname{argmax}_\Theta \frac{1}{Z_\Lambda^\Theta(\gamma_\Lambda^c)} e^{-H_\Lambda^\Theta(\gamma_\Lambda | \gamma_\Lambda^c)}.$$

MLE and MPLE procedures

- P a Quermass Model for $\Theta^* = (z^*, \theta_1^*, \theta_2^*, \theta_3^*)$.
- h^{Θ^*} and h^Θ are the local energies for Θ^* and Θ .
- γ a realization of P .
- **MLE** : In Λ , the MLE is defined by

$$\hat{\Theta}_\Lambda = \operatorname{argmax}_\Theta \frac{1}{Z_\Lambda^\Theta(\gamma_\Lambda^c)} e^{-H_\Lambda^\Theta(\gamma_\Lambda | \gamma_\Lambda^c)}.$$

It is studied by Moller, Helisova (2008) when the intensity z^* is known.

Indeed from the data it is not possible to compute $\#(\gamma_\Lambda)$.

MLE and MPLE procedures

- P a Quermass Model for $\Theta^* = (z^*, \theta_1^*, \theta_2^*, \theta_3^*)$.
- h^{Θ^*} and h^Θ are the local energies for Θ^* and Θ .
- γ a realization of P .
- MLE** : In Λ , the MLE is defined by

$$\hat{\Theta}_\Lambda = \operatorname{argmax}_\Theta \frac{1}{Z_\Lambda^\Theta(\gamma_\Lambda^c)} e^{-H_\Lambda^\Theta(\gamma_\Lambda | \gamma_\Lambda^c)}.$$

It is studied by Moller, Helisova (2008) when the intensity z^* is known.

Indeed from the data it is not possible to compute $\#(\gamma_\Lambda)$.

- MPLE** : In Λ , the MPLE is defined by

$$\hat{\Theta}_\Lambda = \operatorname{argmin}_\Theta \sum_{(x,R) \in \gamma_\Lambda} h^\Theta((x,R), \gamma \setminus (x,R)) + \int e^{-h^\Theta((x,R), \gamma)} \lambda_\Lambda(dx) Q(dR).$$

MLE and MPLE procedures

- P a Quermass Model for $\Theta^* = (z^*, \theta_1^*, \theta_2^*, \theta_3^*)$.
- h^{Θ^*} and h^Θ are the local energies for Θ^* and Θ .
- γ a realization of P .
- MLE** : In Λ , the MLE is defined by

$$\hat{\Theta}_\Lambda = \operatorname{argmax}_\Theta \frac{1}{Z_\Lambda^\Theta(\gamma_{\Lambda^c})} e^{-H_\Lambda^\Theta(\gamma_\Lambda | \gamma_{\Lambda^c})}.$$

It is studied by Moller, Helisova (2008) when the intensity z^* is known.

Indeed from the data it is not possible to compute $\#(\gamma_\Lambda)$.

- MPLE** : In Λ , the MPLE is defined by

$$\hat{\Theta}_\Lambda = \operatorname{argmin}_\Theta \sum_{(x,R) \in \gamma_\Lambda} h^\Theta((x,R), \gamma \setminus (x,R)) + \int e^{-h^\Theta((x,R), \gamma)} \lambda_\Lambda(dx) Q(dR).$$

This procedure does not work since, from the data, we don't know where are the balls.

Takacs-Fiksel procedure.

This procedure is based on the GNZ equilibrium equation :

$$E_P\left(\sum_{X \in \gamma} f(X, \gamma \setminus X)\right) = E_P\left(\int f(X, \gamma) e^{-h^{\Theta^*}(X, \gamma)} z^* \lambda \otimes Q(dX)\right),$$

Takacs-Fiksel procedure.

This procedure is based on the GNZ equilibrium equation :

$$E_P\left(\sum_{X \in \gamma} f(X, \gamma \setminus X)\right) = E_P\left(\int f(X, \gamma) e^{-h^{\Theta^*}(X, \gamma)} z^* \lambda \otimes Q(dX)\right),$$

We have to choose f such that the empirical sum are computable.

Takacs-Fiksel procedure.

This procedure is based on the GNZ equilibrium equation :

$$E_P\left(\sum_{X \in \gamma} f(X, \gamma \setminus X)\right) = E_P\left(\int f(X, \gamma) e^{-h^{\Theta^*}(X, \gamma)} z^* \lambda \otimes Q(dX)\right),$$

We have to choose f such that the empirical sum are computable.

-Examples of such functions f :

- $f_0(X, \gamma) = \mathcal{L}(\partial B(X) \cap \bar{\gamma}^c)$.

In this situation $\sum_{X \in \gamma_\Lambda} f_0(X, \gamma \setminus X) \approx \mathcal{L}(\bar{\gamma}_\Lambda)$.

Takacs-Fiksel procedure.

This procedure is based on the GNZ equilibrium equation :

$$E_P\left(\sum_{X \in \gamma} f(X, \gamma \setminus X)\right) = E_P\left(\int f(X, \gamma) e^{-h^{\Theta^*}(X, \gamma)} z^* \lambda \otimes Q(dX)\right),$$

We have to choose f such that the empirical sum are computable.

-Examples of such functions f :

- $f_0(X, \gamma) = \mathcal{L}(\partial B(X) \cap \bar{\gamma}^c)$.

In this situation $\sum_{X \in \gamma_\Lambda} f_0(X, \gamma \setminus X) \approx \mathcal{L}(\bar{\gamma}_\Lambda)$.

- $f_\alpha(X, \gamma) = \mathcal{L}(\partial B(x, R + \alpha) \cap (\bar{\gamma})_\alpha^c)$.

In this situation $\sum_{X \in \gamma_\Lambda} f_\alpha(X, \gamma \setminus X) \approx \mathcal{L}((\bar{\gamma}_\Lambda)_\alpha)$.

Takacs-Fiksel procedure.

This procedure is based on the GNZ equilibrium equation :

$$E_P\left(\sum_{X \in \gamma} f(X, \gamma \setminus X)\right) = E_P\left(\int f(X, \gamma) e^{-h^{\Theta^*}(X, \gamma)} z^* \lambda \otimes Q(dX)\right),$$

We have to choose f such that the empirical sum are computable.

-Examples of such functions f :

- $f_0(X, \gamma) = \mathcal{L}(\partial B(X) \cap \bar{\gamma}^c)$.

In this situation $\sum_{X \in \gamma_\Lambda} f_0(X, \gamma \setminus X) \approx \mathcal{L}(\bar{\gamma}_\Lambda)$.

- $f_\alpha(X, \gamma) = \mathcal{L}(\partial B(x, R + \alpha) \cap (\bar{\gamma})_\alpha^c)$.

In this situation $\sum_{X \in \gamma_\Lambda} f_\alpha(X, \gamma \setminus X) \approx \mathcal{L}((\bar{\gamma}_\Lambda)_\alpha)$.

- $f_{iso}(X, \gamma) = \mathbb{I}_{B(X) \cap \bar{\gamma} = \emptyset}$.

In this situation $\sum_{X \in \gamma_\Lambda} f_{iso}(X, \gamma \setminus X)$ is equal to the number of isolated balls in $\bar{\gamma}_\Lambda$.

Takacs-Fiksel procedure

For any function f we define

$$\Delta_{f,\Lambda} := \sum_{X \in \gamma_\Lambda} f(X, \gamma \setminus X) - \int f(X, \gamma) e^{-h^\Theta(X, \gamma)} \lambda_\Lambda \otimes Q(dX).$$

Takacs-Fiksel procedure

For any function f we define

$$\Delta_{f,\Lambda} := \sum_{X \in \gamma_\Lambda} f(X, \gamma \setminus X) - \int f(X, \gamma) e^{-h^\Theta(X, \gamma)} \lambda_\Lambda \otimes Q(dX).$$

TFE :

$$\hat{\Theta} := \operatorname{argmin}_\Theta \left(\Delta_{f_1, \Lambda}^2 + \Delta_{f_2, \Lambda}^2 + \Delta_{f_3, \Lambda}^2 + \Delta_{f_4, \Lambda}^2 \right).$$

Takacs-Fiksel procedure

For any function f we define

$$\Delta_{f,\Lambda} := \sum_{X \in \gamma_\Lambda} f(X, \gamma \setminus X) - \int f(X, \gamma) e^{-h^\Theta(X, \gamma)} \lambda_\Lambda \otimes Q(dX).$$

TFE :

$$\hat{\Theta} := \operatorname{argmin}_\Theta \left(\Delta_{f_1, \Lambda}^2 + \Delta_{f_2, \Lambda}^2 + \Delta_{f_3, \Lambda}^2 + \Delta_{f_4, \Lambda}^2 \right).$$

By the GNZ equation :

$$E_P(\Delta_{f_i, \Lambda}) = 0$$

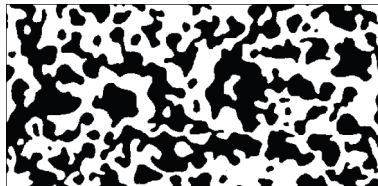
and by ergodicity

$$\frac{1}{|\Lambda|} \Delta_{f_i, \Lambda} \xrightarrow{\Lambda \rightarrow \mathbb{R}^2} 0.$$

References

- The Takacs-Fiksel procedure is introduced in 1984-86 by Takacs and Fiksel
- Asymptotic properties : Consistency and normality under a mixing assumption, Heinrich 92. General case, Coeurjolly, Der., Drouilhet and Lavancier, 2012.
- Application for the Quermass model : Der., Helisova and Lavancier 2013.

Heather Dataset

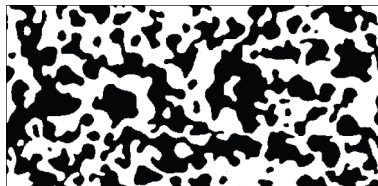


Heather : Real data

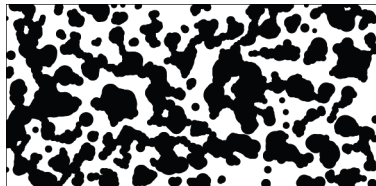


Approximation by balls

Heather Dataset



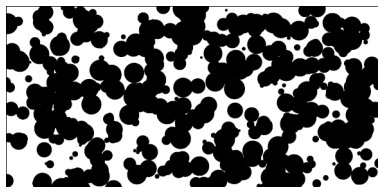
Heather : Real data











Approximation by balls

TFE for Quermass model with Q uniform in $[0, 0.5]$:

$$z = 2.12, \theta_1 = 0, \theta_2 = 0.14 \text{ and } \theta_3 = 0.22.$$



-  A. BADDELEY and M. N. M. VAN LIESHOUT, *Area-interaction point processes*, Ann. Inst. Statist. Math. Vol. 47 no. 4, 601–619 (1995).
-  J.T. CHAYES, L. CHAYES and R. KOTECKY, *The Analysis of the Widom-Rowlinson Model by Stochastic Geometric methods*, Commun. math. Phys. 172, 551-569 (1995).
-  J.-F. COEURJOLLY, Der., R. DROUILHET and F. LAVANCIER, Takacs Fiksel method for stationary marked Gibbs point processes, Scandinavian Journal of Statistics 39(3) (2012) 416-443
-  D. COUPIER and D. Der., Continuum percolation for Quermass model, to appear in Elect. Journal of Prob.
-  D. DER, Existence of Quermass processes for non locally stable interaction and non bounded convex grains, Adv. in Appl. probab. 41 664-681 (2009).
-  D. DER and F. LAVANCIER, Fitting all parameters of the Quermass model by the Takacs-Fiksel method, submitted.
-  T. FIKSEL, *Estimation of interaction potentials of Gibbsian point processes*, Statistics 19, 77-86 (1988).
-  H.-O. GEORGII, O. HÄGGSTRÖM and C. MAES, The random geometry of equilibrium phases, Phase Transitions and Critical Phenomena Volume 18, 2001, Pages 1-142



L. HEINRICH, *Mixing Properties of Gibbsian Point Processes and asymptotic Normality of Takacs-Fiksel Estimates*, Preprint (1992).



W. S. KENDALL, M. N. M. VAN LIESHOUT and A. J. BADDELEY, *Quermass-interaction processes conditions for stability*, Adv. Appl. Prob. 31, 315–342 (1999).



C. N. LIKOS, K. R. MECKE and H. WAGNER, *Statistical morphology of random interface microemulsions*, J. Chem. Phys., 9350–9361 (1995).



K. R. MECKE, *A morphological model for complex fluids*, J. Phys. Condens. Matter 8, 9663–9667 (1996).



J. MØLLER and K. HELISOVA, *Likelihood inference for unions of interacting discs*, to appear in Scandinavian Journal of Statistics (2010).



X.X. NGUYEN and H. ZESSIN, *Integral and differential characterizations of the Gibbs process*, Math. Nach. 88 105-115 (1979).



R. TAKACS, *Estimator for the pair-potential of a Gibbsian point process*, Math. Operationsforsch. u. Statist., ser. statist. 17 428-433 (1986).



B. WIDOM and J. S. ROWLINSON, *New model for the study of liquid-vapor phase transitions*, J. Chem. phys. 52 ,1670–1684 (1970).